| Pressure (kbar) | 0 | 25 | 50 | 75 | 100 | 125 | 150 | 175 |
|--------------------------|-------------------------------|-------|-------|--------------|-------------|-------|-------|-------|
| | | | | Isomer shift | ft (mm/sec) | a | | r |
| Low spin | | | | | | | | |
| $FePc(Py)_2$ | 0.270 | 0.245 | 0.200 | 0.175 | 0.167 | 0.165 | 0.162 | 0.160 |
| FePc(3-Pic) ₂ | 0.260 | 0.240 | 0.200 | 0.175 | 0.167 | 0.165 | 0.162 | 0.160 |
| FePc(4-Pic) ₂ | 0.270 | 0.265 | 0.215 | 0.180 | 0.167 | 0.165 | 0.162 | 0.160 |
| FePc(Pip) ₂ | 0.265 | 0.255 | 0.210 | 0.166 | 0.140 | 0.132 | 0.128 | 0.125 |
| Intermediate spin | | | | | | | | |
| FePc | 0.385 | 0.373 | 0.350 | 0.324 | 0.302 | 0.290 | 0.280 | 0.274 |
| FePc(Py) ₂ | | | 0.335 | 0.295 | 0.280 | 0.270 | 0.260 | 0.255 |
| FePc(3-Pic) ₂ | | | 0.330 | 0.295 | 0.273 | 0.266 | 0.260 | 0.255 |
| FePc(4-Pic) ₂ | | | 0.322 | 0.300 | 0.285 | 0.275 | 0.275 | 0.270 |
| $FePc(Pip)_2$ | | ••• | ••• | 0.245 | 0.255 | 0.252 | 0.245 | 0.240 |
| | Quadrupole splitting (mm/sec) | | | | | | | |
| Low spin | | | | | | | | |
| $FePc(Py)_2$ | 2.01 | 1.97 | 1.93 | 1.92 | 1.93 | 1.92 | 1.91 | 1.89 |
| FePc(3-Pic) ₂ | 1.93 | 1.92 | 1.89 | 1.89 | 1.91 | 1.93 | 1.93 | 1.92 |
| FePc(4-Pic) ₂ | 2.03 | 2.03 | 1.97 | 1.95 | 1.93 | 1.92 | 1.91 | 1.89 |
| FePc(Pip) ₂ | 2.22 | 2.26 | 2.23 | 2.16 | 2.10 | 2.08 | 2.07 | 2.07 |
| Intermediate spin | | | | | | | | |
| FePc | 2.58 | 2.74 | 2.90 | 3.06 | 3.14 | 3.18 | 3.21 | 3.23 |
| $FePc(Py)_2$ | | | 2.78 | 2.95 | 3.02 | 3.05 | 3.07 | 3.08 |
| FePc(3-Pic) ₂ | | | 2.76 | 2.95 | 3.04 | 3.09 | 3.12 | 3.15 |
| FePc(4-Pic) ₂ | | | 2.78 | 2.96 | 3.04 | 3.08 | 3.10 | 3.12 |
| FePc(Pip), | | | | 2.72 | 2.83 | 2.93 | 3.01 | 3.04 |

TABLE II. Mössbauer parameters for FePc(L)₂.

^a Relative to iron metal.

it is possible to measure only order of magnitude changes in absolute intensity, so it is impossible to determine if the absolute intensity of the N transition is changing with pressure.

It has been established that the ground state of FePc is $3_{B_{2g}}$.¹² The order of the *d* orbitals and their symmetries in D_{4h} are: d_{xz} , $d_{yz}(e_g)$; $d_{xy}(b_{2g})$; $d_{z^2}(a_{1g})$; $d_{x^2-y^2}(b_{1g})$. This order is slightly different from the order that Zerner et al.¹⁹ calculated for the ferrous porphyrin complex. Their energy scheme had the d_{xy} orbital lower than the d_{xz} , $y_z(d_\pi)$ orbitals. Their calculations indicated that the metal d_{π} orbitals were approximately 90% pure, indicating that there is some backbonding of the metal d_{π} orbitals to the porphyrin antibonding orbitals. Since the lowest-lying ligand antibonding orbital for both the porphyrins and the phthalocyanines is an $e_q(\pi)$ orbital, backbonding of metal d_{π} electrons to this level is the most likely case. The additional resonance structure of FePc should allow for even more backbonding and thus a greater stabilization of the metal d_{π} orbitals. Therefore, in FePc the metal e_{q} orbital lies lower in energy than the metal b_{1g} orbital.

When additional ligands are coordinated perpendicular to the phthalocyanine plane, two additional perturbations of the energies of the metal d orbitals take place. The first is that the d_z^2 orbital is raised in energy due to bonding of the axial ligands. This increase in energy is greater than the sipn pairing energy, so that the electron formerly in the d_z^2 orbital is paired with the electron in the d_{xy} orbital, forming the low spin complex. Also, if the axial ligands have empty π^* orbitals, the d_{π} electrons can backbond to those orbitals, causing a further decrease in the energy of the d_{π} orbital. The effect of this additional backbonding can be shown by the energy of the charge transfer transition in these low spin complexes. The more backbonded the metal π orbitals are, the higher in energy the charge transfer transition will be. According to Dale,9 the order of increasing charge transfer energy is: n-butylamine, NH₃<piperdine<imidazole<pyridine. This is also the order of increasing π acceptor character of these ligands, with the first three having comparable ability. It was found that the picoline adducts studied have lower charge transfer energies than the pyridine adduct.

In the piperidine adduct the axial ligands have no empty π^* orbitals available for backbonding. The fact that it is low spin is due to the very strong σ bonding which raises the energy of the d_y^2 orbital. This is exemplified by the fact that the acid dissociation constant pK_a for the piperidine derivative is 11.2 compared with 5.25, 5.60, and 6.0 for the pyridine, 3-picoline, and 4-picoline adducts.²⁰

In the previous section it was shown that the low spin pyridine, picoline, and piperidine adducts of FePc convert to a different spin state with increasing pressure.

The major cause for this conversion would seem to be a reduction in the backbonding of metal electrons to the empty π^* orbitals of the ligand. As indicated earlier, the $\pi - \pi^*$ transitions all shift to lower energy with increasing pressure. The effect of this is to move charge into the e_q^* ligand antibonding orbitals. The shift of the lowest $\pi - \pi^*$ peak is approximately 10% of the total energy of the peak maximum. However, the electron transfer is a thermal process. Figure 12 shows a schematic configuration coordinate diagram for the $\pi - \pi^*$ system. A modest decrease in energy of the optical transition can be sufficient to cause a thermal population of the excited state as shown since the thermal transition is not restricted by the Franck-Condon principle. The e_q^* orbitals are the same ones that are available to the d_{π} electrons for backbonding. The effect of this additional ligand π electron character will be to decrease the backbonding of the metal d_{π} electrons to them. This in turn will raise the energy of the d_{π} orbitals and decrease the energy difference between the d_z^2 and the d_{π} orbitals sufficiently to cause a change in the spin state of the ferrous iron with pressure. A transition to higher spin state has been observed in some ferrocyanides²¹ and iron phenanthroline compounds,22 where it has also been attributed to decreasing backbonding.

At high pressure the conversion to the new spin state stops. This appears to represent a balance between the reduced backbonding which tends to produce a higher spin state and then the delocalization of the 3d electrons at high pressure which favors a lower spin state. The delocalization is facilitated by the thermal transfer of the ligand electron to the π^* orbitals which are nearer the perimeter of the molecule leaving room for greater expansion of the 3d orbitals.



MIXING OF LIGAND $\pi - \pi^*$ LEVELS AT HIGH PRESSURE





FIG. 13. Mössbauer spectra of sheared FePc.

The picoline adducts converted more to intermediaet spin than did the pyridine adduct. The methyl substitution on the pyridine decreases the ability to accept metal d_{π} electrons. This is reflected by the lower energy of the charge transfer band. Since the energy difference between the d_{π} and d_z^2 orbitals is then less in the picoline adducts, their increased conversion to intermediate spin is expected. The piperidine derivative showed the least conversion to intermediate spin. In this compound the only backbonding is to the phthalocyanine, so that the decrease in backbonding is less than for the other adducts, and the conversion to intermediate spin is less. The conversion actually decreases at the higher pressures where the spreading of the 3d orbitals more than compensates for the decrease in backbonding.²³

It is possible that changes in π bonding and σ bonding also contribute to the spin change, but we believe these effects are secondary.

The new spin state has been labeled intermediate spin because its isomer shift and quadrupole splitting are very nearly those of FePc. To obtain an intermediate spin state it is necessary to remove the degeneracy of the d_{z^2} and $d_{x^2y^2}$ orbitals by a significant amount. It is perhaps surprising that this occurs in the pseudooctahedral symmetry of the pththalocyanine deriva-